

## Electromagnetic Theory

### (A) Maxwell's field equations:

E.M. theory is a subject which deals with electrostatic as well as magnetostatic fields. A static charge can produce electrostatic field whereas a charge in motion can produce magnetostatic field and e.m. theory can be renamed as electrodynamics. An e.m. wave can be characterized by electric field ( $\vec{E}$ ), magnetic field ( $\vec{H}$ ), electric displacement vector ( $\vec{D}$ ) and magnetic induction ( $\vec{B}$ ). These field vectors obey four fundamental equations which are known as Maxwell's field equations. These equations are not absolutely new but extraction from and modification of existing laws in electrostatics and magnetostatics. Let us now derive them.

### (a) Maxwell's 1st equation $[\vec{\nabla} \cdot \vec{D} = \rho]$ :-

Derivation: we have Gauss's theorem in electrostatics which states that "flux of electric field across any closed surface is equal to  $1/\epsilon_0$  times the charge enclosed by the surface".

Mathematically, 
$$\oiint \vec{E} \cdot \hat{n} ds = \frac{1}{\epsilon_0} \rho$$

or, 
$$\oiint \vec{E} \cdot \hat{n} ds = \frac{1}{\epsilon_0} \left[ \iiint \rho dv - \iiint \vec{\nabla} \cdot \vec{P} dv \right]$$

where  $\rho =$  free vol. charge density

$\rho_p = -\vec{\nabla} \cdot \vec{P} =$  Polarization charge density.



$\vec{P}$  = Polarization having unit  $C/m^2$ .

Applying Gauss's divergence theorem in LHS and Rearranging,

$$\int_V (\vec{\nabla} \cdot \epsilon_0 \vec{E} + \vec{\nabla} \cdot \vec{P}) dv = \iiint \rho dv.$$

a,  $\iiint_V \vec{\nabla} \cdot (\epsilon_0 \vec{E} + \vec{P}) dv = \iiint \rho dv.$

or,  $\iiint (\vec{\nabla} \cdot \vec{D} - \rho) dv = 0$ , where  $\vec{D} = \epsilon_0 \vec{E} + \vec{P}$   
= electric disp. vector

Since elementary volume  $dv$ , is taken arbitrarily, so it can be <sup>not</sup> zero. Hence above eqn. holds good if and only if

$$\vec{\nabla} \cdot \vec{D} = \rho$$

①

Phy. Sig

This is Maxwell's 1st equation.

It is nothing but differential form of Gauss's theorem in electrostatics which states that " " "

(b) Maxwell's 2nd equation  $[\vec{\nabla} \cdot \vec{B} = 0]$

Derivation: Since isolated magnetic poles does not exist, according to Gauss's theorem in magnetostatics, we can write, flux of mag. field over any closed surface is zero i.e.

$$\oint \vec{B} \cdot \hat{n} ds = 0$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

Again,  $\iiint \vec{\nabla} \cdot \vec{B} dv = 0$  since magnetic forces are closed curves / lines of goes of to

infinitesimal, we can say, magnetic ind. is a flux free/solenoidal vector; i.e. we can write

$$\oint \vec{\nabla} \cdot \vec{B} = 0$$

This is Maxwell's 2nd eqn.  
Physical Significance; HW

③ Maxwell's 3rd equation;  $[\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}]$

We have Faraday's law of e.m. induction; i.e. induced emf = rate of change of mag. flux  

$$e = -\frac{\partial \psi}{\partial t} \quad (4)$$

now emf  $e =$  work done to bring unit positive charge once round a closed path in the el. field

$$= \oint_c \vec{F} \cdot d\vec{r} = \oint_c \vec{E} \cdot d\vec{r}$$

and  $\psi = \iint_s \vec{B} \cdot \hat{n} ds$

$$\therefore \oint_c \vec{E} \cdot d\vec{r} = -\frac{\partial}{\partial t} \iint_s \vec{B} \cdot \hat{n} ds$$

App Stokes's Theorem in RHS,

$$\iint_s (\vec{\nabla} \times \vec{E}) \cdot \hat{n} ds = -\iint_s \frac{\partial \vec{B}}{\partial t} \cdot \hat{n} ds$$

$$\therefore \iint_s \left( \vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} \right) \cdot \hat{n} ds = 0$$

$$\boxed{\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}} \quad \text{--- (3)}$$

This is Maxwell's 3rd eqn. which is nothing but diff. form of Faraday's law of induction.



(a) Maxwell's 4th equation:  $\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$   
 [Copy from Part-I]

(B) Equation of continuity  $\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$

(c) Home work: Copy " from Part-I

Constitutive relations: The following relations are known as Constitutive relations:

(i)  $\vec{D} = \epsilon \cdot \vec{E}$ ,  $\left\{ \begin{array}{l} \epsilon = \text{Permittivity of the medium} \\ = \epsilon_0 \epsilon_r \end{array} \right.$

(ii)  $\vec{B} = \mu \vec{H}$ ,  $\left\{ \begin{array}{l} \epsilon_0 = \text{Per. of free space} \\ = 8.854 \times 10^{-12} \text{ F/m} \\ \epsilon_r = \text{relative Per./SVC/} \\ \text{dielectric constant} \end{array} \right.$

$\mu = \text{Permeability}$   
 $= \mu_0 \mu_r$

$\mu_0 = 4\pi \times 10^{-7} \text{ H/m}$ ,  $\left[ \frac{1}{4\pi\epsilon_0} \right] = 9 \times 10^9 \text{ m/F}$

(iii) ohm's law:  $\vec{J} = \sigma \vec{E}$  (H.W / Prove it)

Maxwell's equations for diff. media:

We have Maxwell's field equations:

$$\vec{\nabla} \cdot \vec{D} = \rho$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}, \quad \vec{J} = \sigma \vec{E}$$

(a) Free Space: Free space is characterized by  $\rho = 0, \sigma = 0$

$\epsilon = \epsilon_0, \mu = \mu_0$   
 So field equations become:

$$\vec{\nabla} \cdot \vec{E} = 0 \quad \text{--- (1)}$$

$$\vec{\nabla} \cdot (\mu_0 \vec{H}) = 0$$

$$\text{or } \vec{\nabla} \cdot \vec{H} = 0 \quad \text{--- (2)}$$

$$\vec{\nabla} \times \vec{E} = -\mu_0 \frac{\partial \vec{H}}{\partial t}$$

$$\vec{\nabla} \times \vec{H} = \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

(b) Dielectric medium. (c) Conducting medium

v.u/94 i) We have Maxwell's 3rd eqn, i.e. diff. form of Faraday's law;

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Taking divergence on both sides:

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{E}) = -\vec{\nabla} \cdot \frac{\partial \vec{B}}{\partial t}$$

$$\text{or, } 0 = -\frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{B})$$

Since  $\vec{\nabla} \cdot \vec{B}$  is not independent of time, above eqn. holds good, if  $\vec{\nabla} \cdot \vec{B} = 0$

} Since differentiation w.r.t. time and space are interchangeable.

This is Gauss's theorem in mag.

i.e. 2nd law.

ii)

4th  $\rightarrow$  1st

We have modified Ampere's law;



$$\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

Taking div. on both sides:

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{H}) = \vec{\nabla} \cdot \vec{J} + \vec{\nabla} \cdot \frac{\partial \vec{D}}{\partial t}$$

or,  $0 = -\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{D})$

[Since, ... ,  $\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$ ]

$$\frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{D} - \rho) = 0$$

$$\vec{\nabla} \cdot \vec{D} = \rho$$

Since fields are time varying.

This is Gauss's law in electrostatics.

V.U/95

8. (a)

We have Maxwell's 4th eqn:

$$\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

Taking div. on both sides:

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{H}) = \vec{\nabla} \cdot \vec{J} + \vec{\nabla} \cdot \frac{\partial \vec{D}}{\partial t}$$

$$0 = \vec{\nabla} \cdot \vec{J} + \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{D}), \text{ since diff.}$$

From Max. 1st eqn:  $\vec{\nabla} \cdot \vec{D} = \rho$

$$\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$$